

## THE SECOND LAW OF THERMODYNAMICS

According to the First Law of Thermodynamics, the energy of the universe remains constant in any change. Energy may be transformed from one form into another, but the energy of the universe cannot be changed.

However, all real changes have a natural direction. The transformation in the opposite direction is unnatural. In nature, rivers run from the mountains to the sea, never in the opposite way. An insulated metal rod initially hot at one end and cold at the other comes to a uniform temperature. The first law does not give us any information regarding the natural direction. As long as the energy of the universe remains constant, the change may occur in either direction and satisfy the first law. The property that indicates the natural direction is "**entropy**" which is defined by the Second Law of Thermodynamics.

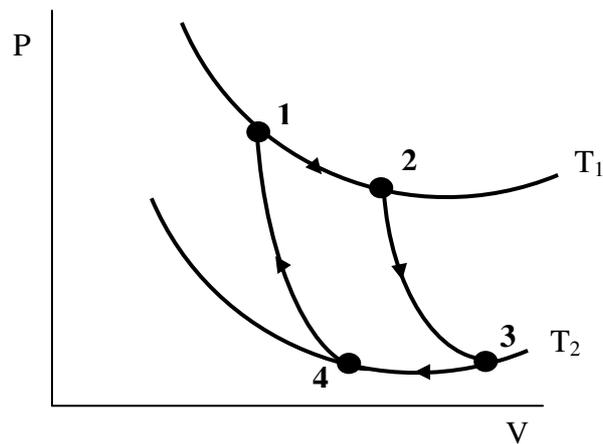
### The Second Law of Thermodynamics

The development of the Second Law is based on the studies on heat engines. A **heat engine** transforms thermal energy into mechanical energy. The Second Law states that it is impossible for an engine to extract heat from the surroundings and completely transform this energy to mechanical energy. The principals governing the transformation of thermal energy into mechanical energy have been investigated by French engineer Sadi Carnot. All these principals are based on a reversible cycle called "**Carnot Cycle**".

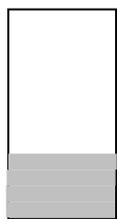
### Carnot Cycle

If there are two heat reservoirs at different temperatures, thermal energy may partially be transformed into mechanical energy by a cyclic operation of a fluid, when

heat is flowing from one reservoir to the other. Since a cyclic operation brings the system back to its initial condition,  $\Delta E = \Delta H = 0$ .



The Carnot Cycle consists of 4 reversible steps;



$T_1, P_1, V_1$



$T_1, P_2, V_2$



$T_2, P_3, V_3$



$T_2, P_4, V_4$

- 1. Reversible Isothermal Expansion:** The gas is confined in a cylinder. The cylinder is immersed in a heat reservoir at  $T_1$  and expands isothermally from  $V_1$  to  $V_2$ .

$$(T_1, P_1, V_1) \rightarrow (T_1, P_2, V_2)$$

$$\Delta E_1 = Q_1 - W_1$$

- 2. Reversible Adiabatic Expansion:** The cylinder is taken out of the reservoir, insulated and the gas expands adiabatically from  $V_2$  to  $V_3$ . In this step, the temperature of the gas drops from  $T_1$  to  $T_2$ .

$$(T_1, P_2, V_2) \rightarrow (T_2, P_3, V_3)$$

$$\Delta E_2 = -W_2$$

3. **Reversible Isothermal Compression:** The insulation is removed and the cylinder is immersed into a heat reservoir at  $T_2$ . The gas is compressed isothermally from  $V_3$  to  $V_4$  at.

$$(T_2, P_3, V_3) \rightarrow (T_2, P_4, V_4)$$

$$\Delta E_3 = Q_2 - W_3$$

4. **Reversible Adiabatic Compression:** The cylinder is taken out of the reservoir, insulated and the gas is compressed adiabatically to the original state. In this step, temperature rises from  $T_2$  to  $T_1$ .

$$(T_2, P_4, V_4) \rightarrow (T_1, P_1, V_1)$$

$$\Delta E_4 = -W_4$$

For the cycle;

$$\Delta E = \Delta E_1 + \Delta E_2 + \Delta E_3 + \Delta E_4 = 0$$

$$Q_1 + Q_2 = W_1 + W_2 + W_3 + W_4$$

$$\Downarrow$$

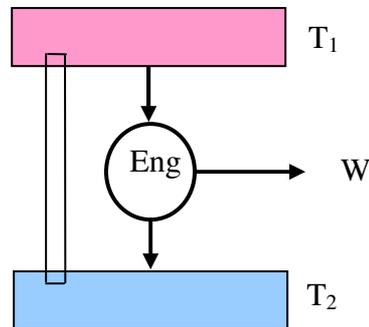
$$Q_{\text{cyc}}$$

$$\Downarrow$$

$$W_{\text{cyc}}$$

Carnot engine operates according to this cycle. It is the simplest heat engine. In order to obtain work from the cycle there must be two heat reservoirs. It is impossible for a system operating in a cycle and connected to a single heat reservoir to produce work in the surroundings.

Carnot engine can be shown schematically as in the figure. According to the Second Law  $Q_1$  and  $Q_2$  cannot have the same sign. If they are both positive, then



$W_{\text{cyc}}$  is also positive. If  $Q_2$  is positive, heat flows out of the cold reservoir. We can restore this quantity by connecting the two reservoirs with a metal rod. Heat flows from the high-temperature reservoir to the low-temperature reservoir. Finally, the reservoir at  $T_2$  gains the energy it lost. Thus, heat engine and low-temperature reservoir behave as a composite engine connected to a single reservoir, which is impossible. So,  $Q_1$  and  $Q_2$  must differ in sign.

The **efficiency** of a heat engine is defined as the ratio of the work produced to the quantity of heat extracted from the high-temperature reservoir.

$$\varepsilon = \frac{W}{Q_1} = \frac{Q_1 + Q_2}{Q_1} = 1 + \frac{Q_2}{Q_1}$$

- The efficiency of any engine is less than or equal to the efficiency of a reversible engine, both operating between the same two heat reservoirs.
- All reversible engines operating between the same two heat reservoirs have the same efficiency.
- The efficiency of an irreversible engine is less than the reversible engine.

## Carnot Cycle of an Ideal Gas

### 1. Reversible Isothermal Expansion

$$\Delta E_1 = Q_1 - W_1 = Q_1 - RT_1 \ln \frac{V_2}{V_1} = 0$$

### 2. Reversible Adiabatic Expansion

$$\Delta E_2 = -W_2 = \int_{T_1}^{T_2} C_v dT$$

### 3. Reversible Isothermal Compression

$$\Delta E_3 = Q_2 - W_3 = Q_2 - RT_2 \ln \frac{V_4}{V_3}$$

### 4. Reversible Adiabatic Compression

$$\Delta E_4 = -W_4 = \int_{T_2}^{T_1} C_v dT$$

Total work obtained from the cycle is;

$$W = RT_1 \ln \frac{V_2}{V_1} - \int_{T_1}^{T_2} C_v dT + RT_2 \ln \frac{V_4}{V_3} - \int_{T_2}^{T_1} C_v dT$$

Interchanging the limits of the second integral we obtain;

$$W = RT_1 \ln \frac{V_2}{V_1} + RT_2 \ln \frac{V_4}{V_3}$$

Since  $V_2 \rightarrow V_3$  and  $V_4 \rightarrow V_1$  are adiabatic changes;

$$T_1 V_2^{\gamma-1} = T_2 V_3^{\gamma-1}$$

$$T_1 V_1^{\gamma-1} = T_2 V_4^{\gamma-1}$$

Dividing both sides we obtain;

$$\frac{V_2}{V_1} = \frac{V_3}{V_4}$$

By substituting this equation in the equation giving  $W$ , we obtain a nice equation for the work obtained from the cycle;

$$W = R(T_1 - T_2) \ln \frac{V_2}{V_1}$$

Work produced by the cycle depends upon the difference between the temperatures of the heat reservoirs and  $V_2/V_1$  ratio. By using the first law equation for step 1 we may write  $Q_1$  as;

$$Q_1 = RT_1 \ln \frac{V_2}{V_1}$$

The efficiency of the cycle may be found as;

$$\varepsilon = \frac{W}{Q_1} = \frac{R(T_1 - T_2) \ln \frac{V_2}{V_1}}{RT_1 \ln \frac{V_2}{V_1}} = \frac{T_1 - T_2}{T_1} = 1 - \frac{T_2}{T_1}$$

The efficiency depends upon the temperature difference.

## Entropy

The efficiency of the Carnot engine may be shown by the below two equations;

$$\varepsilon = 1 + \frac{Q_2}{Q_1} \quad \text{and} \quad \varepsilon = 1 - \frac{T_2}{T_1}$$

---

Subtracting these two equations yields;

$$\frac{Q_1}{T_1} + \frac{Q_2}{T_2} = 0$$

The left-hand side of this equation is the sum of the quantity  $Q/T$  for the cycle.

$$\oint \frac{dQ}{T} = 0$$

Since cyclic integral is zero,  $dQ/T$  is a state function. The defining equation for entropy is;

$$dS = \frac{dQ_{rev}}{T}$$

Entropy is a state function.  $dS$  is an exact differential. The change in entropy does not depend upon the path. Entropy is an extensive property. For a reversible change; the change in entropy may be expressed as;

$$\Delta S = S_2 - S_1 = \int_1^2 \frac{dQ}{T}$$

For all reversible cycles;

$$\oint \frac{dQ}{T} = 0$$

For all irreversible cycles;

$$\oint \frac{dQ}{T} < 0$$

### Clausius Inequality

Assume that the system is transformed irreversibly from 1 to 2, and then brought back reversibly from 2 to 1. For the cycle;

$$\oint \frac{dQ}{T} = \int_1^2 \frac{dQ_{irr}}{T} + \int_2^1 \frac{dQ_{rev}}{T} < 0$$

$$\int_1^2 \frac{dQ_{irr}}{T} + \int_2^1 dS < 0 \qquad \int_1^2 dS \geq \int_1^2 \frac{dQ}{T}$$

This inequality is known as "Clausius Inequality" and simply expressed as;

$$dS \geq \frac{dQ}{T}$$

Clausius inequality is valid for all changes. If the system is isolated  $dQ = 0$ . Then the entropy change becomes;

$$dS \geq 0$$

Any natural change occurring in an isolated system is attended by an increase in entropy. The entropy continues to increase as long as the change occurs. Finally, the system reaches equilibrium. At equilibrium, the entropy is at its maximum value.

Entropy is a measure of disorder. Therefore, increase in entropy means increase in disorder. In the universe, the entropy increases continuously as natural changes occur within it.

Entropy is related with the probability of occurrence. If entropy increases during a change, the probability of occurrence for that change is high.